

FROM RESEARCH TO INDUSTRY



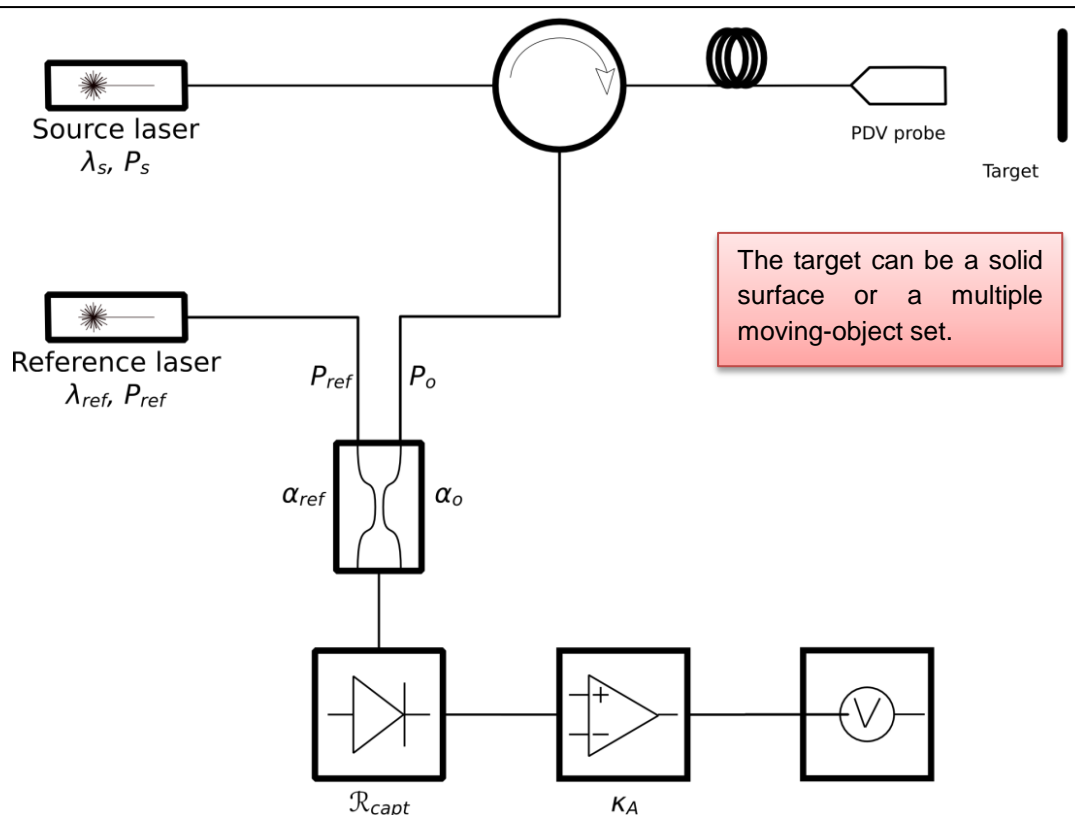
PDV: NOISE ANALYSIS USING SHORT-TIME FOURIER TRANSFORM (STFT)

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PDV User Workshop

LIVERMORE, JUNE, 2016

INTRODUCTION

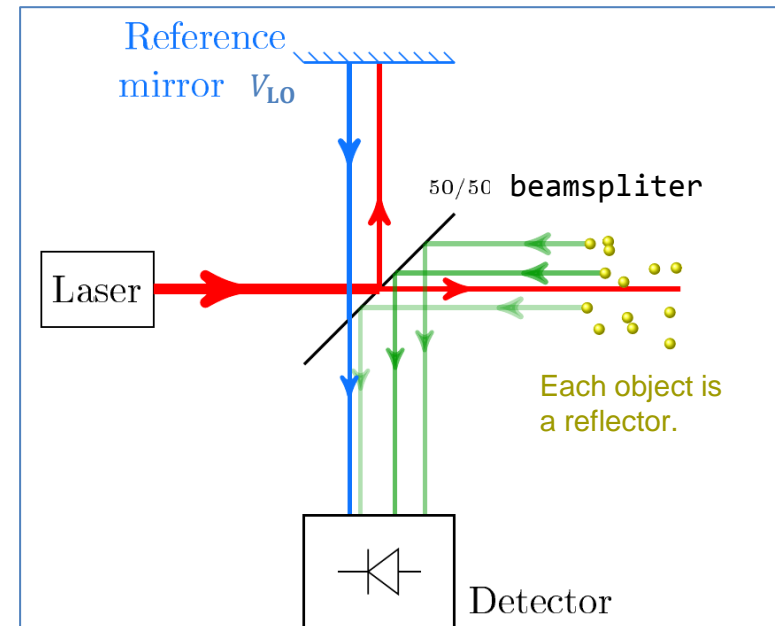
WE CONSIDER THE FOLLOWING PDV DESIGN:



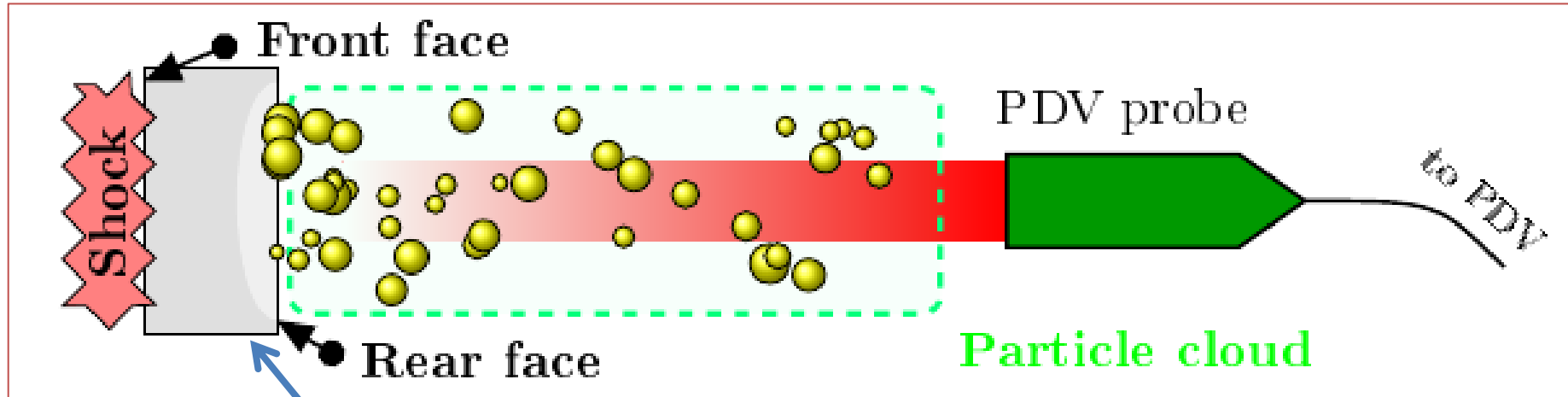
The design is **equivalent to a Michelson interferometer**.

(It is equivalent to have a reference mirror moving with a constant velocity V_{LO} (equal to $\lambda_s/2 \Delta f$))

- Working @ $1.55 \mu\text{m}$,
- Two laser configuration: source and local oscillator.

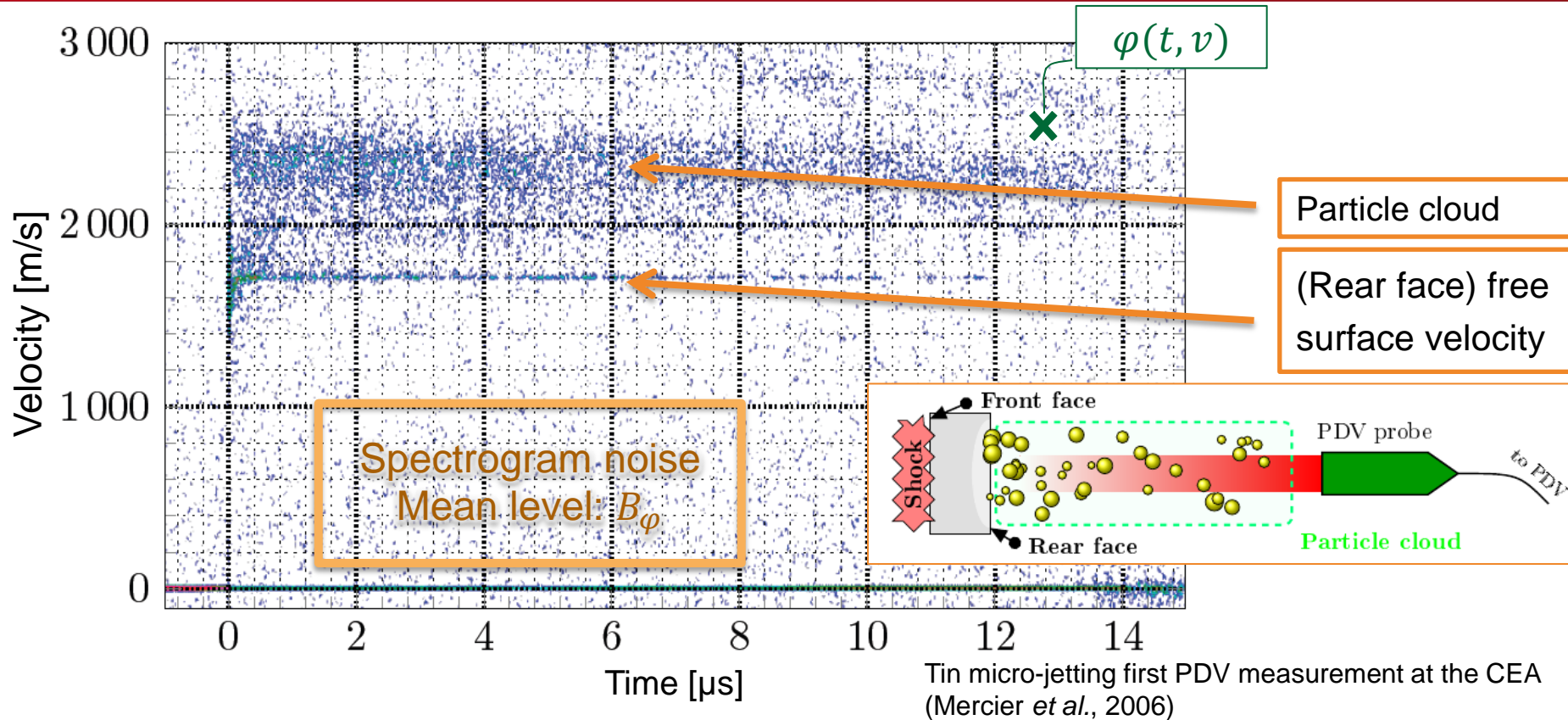


HIGH-VELOCITY PARTICLE CLOUDS IN SHOCK PHYSICS



- A [metallic (tin) plate] undergoes a shockwave (produced by a laser shock, high-explosives or impact)
- Many mechanisms (spalling, microspalling, microjetting, melting, ...) generate a **cloud of micron-sized particles accelerating up to several km/s.**

THE NEED: PARTICLE CLOUDS



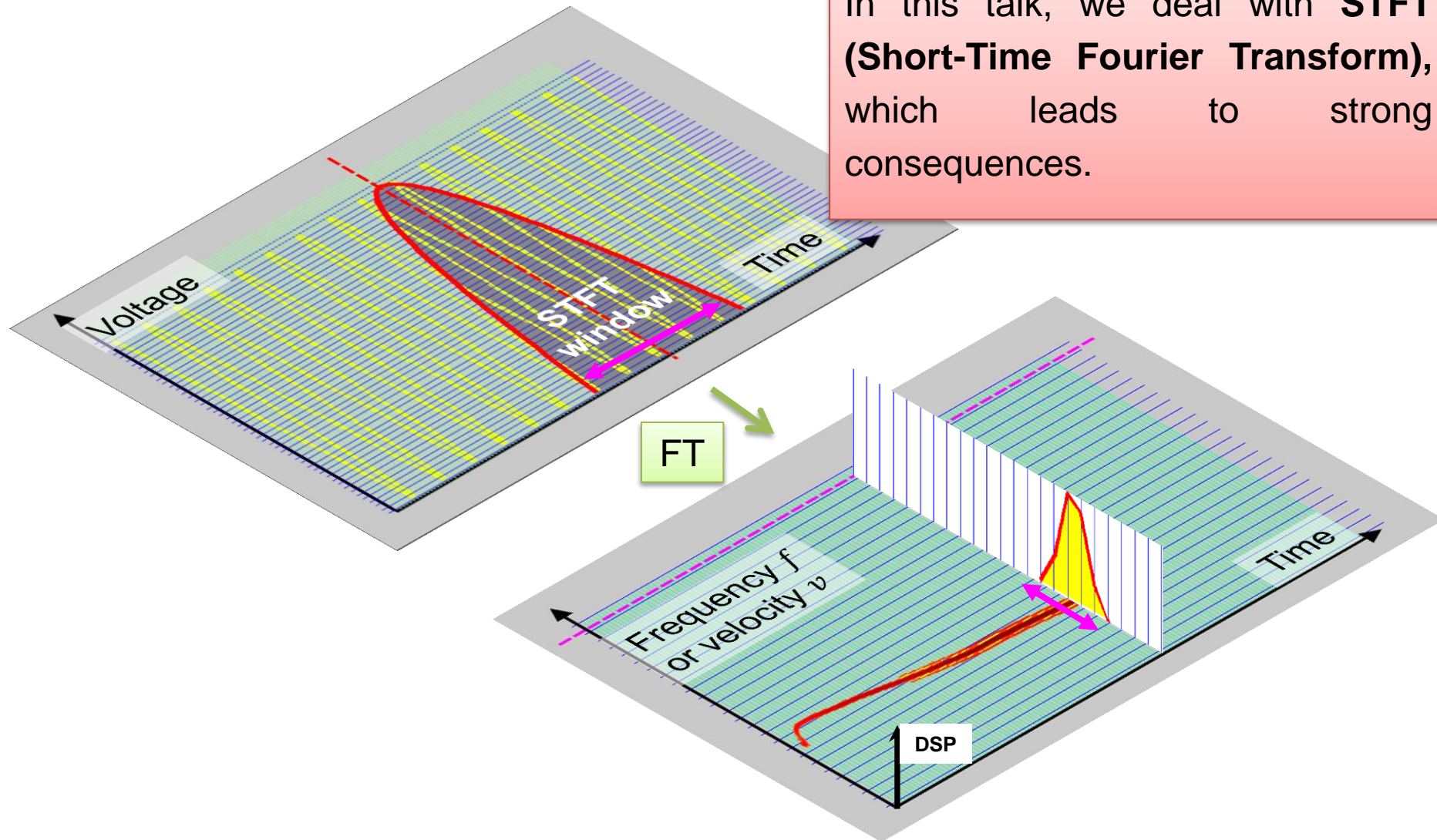
Are the fastest particles detected?

How deep is the cloud analyzed?

- Noise on spectrogram may limit the visibility of the cloud by PDV.
- Noise is linked to the PDV design.

SHORT-TIME FOURIER TRANSFORM

In this talk, we deal with **STFT (Short-Time Fourier Transform)**, which leads to strong consequences.



**WHAT ARE THE NOISE SOURCES OF
THE STANDARD PDV?**

(STANDARD = NOT MPDV)

HOW DO WE SEPARATE THE NOISE SOURCES?

(Ph): shot noise

- Quantum optics : the photons are inseparable amounts of power.

(Q): digitizer noise

- Generally expressed as a number of “equivalent bit”.

(D): Dark noise

Includes background IR (BLIR) noise, electronic noise, thermal noise, ...

(S): spontaneous-emitted photons – ASE noise

- Due to the lasers and amplifiers of the PDV line.

- We will express all noise sources as an optical “NEP” (Noise Equivalent Power”, unit : W/\sqrt{Hz}).**
 - This choice will be explained in the 2nd part.
 - Relation between optical power P_{det} and electrical u_e : $u_e = \mathcal{R} P_{det}$.
- First, we derive the power P_{det} on the PDV detector in order to estimate noise levels.**

BACK TO BASICS: THE OPTICAL POWER COLLECTED BY THE DETECTOR

- For a single free surface: $P_{det}(t) = \alpha_{LO}P_{LO} + \sqrt{\alpha_{LO}P_{LO}\alpha_oP_p} \kappa_{LO}^p \cos[\Omega_p t + \phi_{LO}^p] + \alpha_o P_p$
- For N objects (i.e, a particle cloud), in order of importance:

$$P_{det}(t) = \underbrace{\alpha_{LO}P_{LO}}_{(1)} + \underbrace{\sum_p \sqrt{\alpha_{LO}P_{LO}\alpha_oP_p} \kappa_{LO}^p \cos[\Omega_p t + \phi_{LO}^p]}_{(2)} + \underbrace{\alpha_o \sum_{p,q} \sqrt{P_p P_q} \kappa_p^q \cos[(\Omega_p - \Omega_q)t + \phi_p^q]}_{(3)}$$

- p (or q) : points to a reflector (i.e. a free surface or a particle),
- α_{LO}, α_o : coupling efficiencies,
- κ_p^q : polarization dot product (between p and q),
- $\Omega_p = \frac{4\pi}{\lambda_{PDV}}(v_p - V_{LO})$: relative angular frequency.

- (1) **constant term**, **strong but useless**.
 - \approx few 100 μW ?
- (2) **beating term** (LO coupling), **weak but that's we are looking for!**
 - $\approx \sqrt{(\text{few } 100 \mu W)(? \text{ to } 1 \mu W)} \lesssim 10 \mu W$ for each p .
 - $N \times (\text{few } 10 \mu W)$ spread over the velocity spectrum.
- (3) **cross-beating term**, **very weak, neglected in general!**
 - $N^2 \times (\lesssim 1 \mu W)$: weak when N is not too large.
 - the sum could be separated in two parts:
 - (3c) $p = q, \alpha_o \sum_p P_p$: **constant part**.
 - (3n) $p \neq q, \alpha_o \sum_{(p \neq q)} X_p^q \cos[(\Omega_p - \Omega_q)t + \phi_p^q]$: (nearly) random freq. -> **noise ?!**

WHICH TERMS ARE IMPORTANT FOR THE NOISE? 1/4

Shot noise (or photonic noise)

Because a lot of beats is recorded inside the FT window $\left(T \gg \frac{2\pi}{\Omega_p}\right)$, we consider the constant part of the signal:

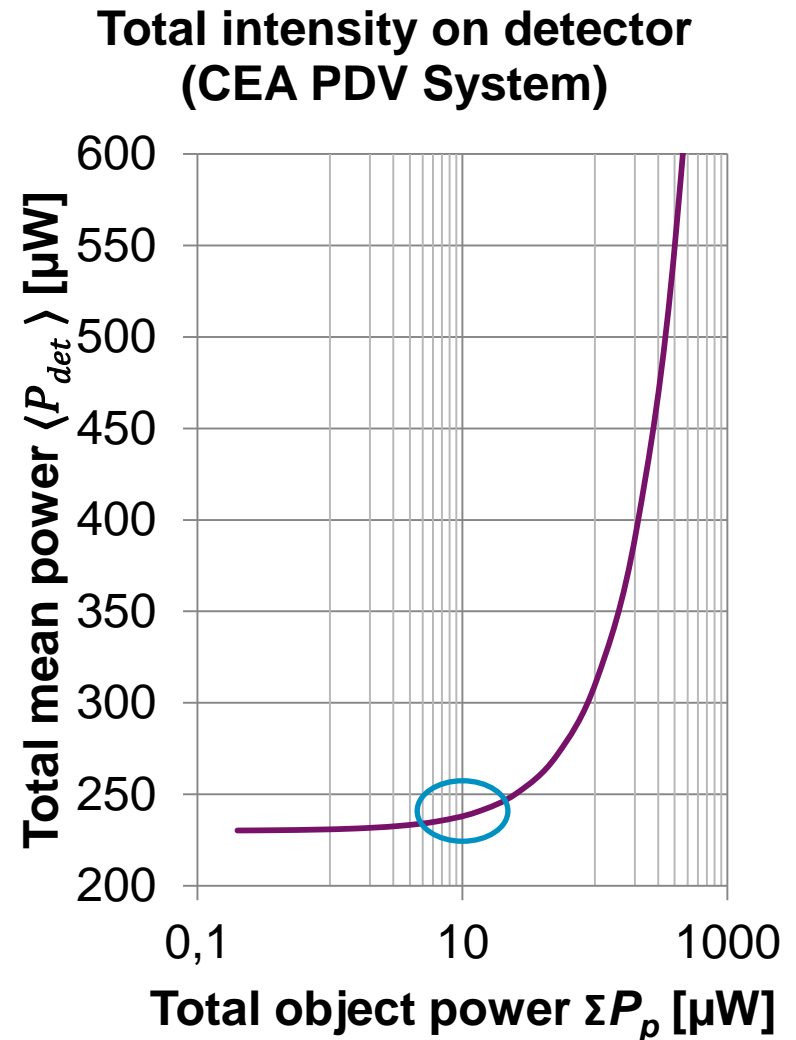
$$\langle P_{det} \rangle = \underbrace{\alpha_{LO} P_{LO}}_{(1)} + \underbrace{\alpha_o \sum_p P_p}_{(3c)}$$

If $\alpha_o \sum_p P_p \leq 10 \mu\text{W}$, we suppose that P_{det} is constant and equal to (1).

The shot-noise spectrum density is:

$$\langle P_{ph} \rangle = K_{ph} \sqrt{f_{BW}} = \sqrt{\frac{2 h c_0 f_{BW} \alpha_{LO} P_{LO}}{\eta \lambda_{PDV}}}$$

$$K_{ph} = \sqrt{\frac{2 h c_0 \alpha_{LO} P_{LO}}{\eta \lambda_{PDV}}}$$



WHICH TERMS ARE IMPORTANT FOR THE NOISE? 2/4

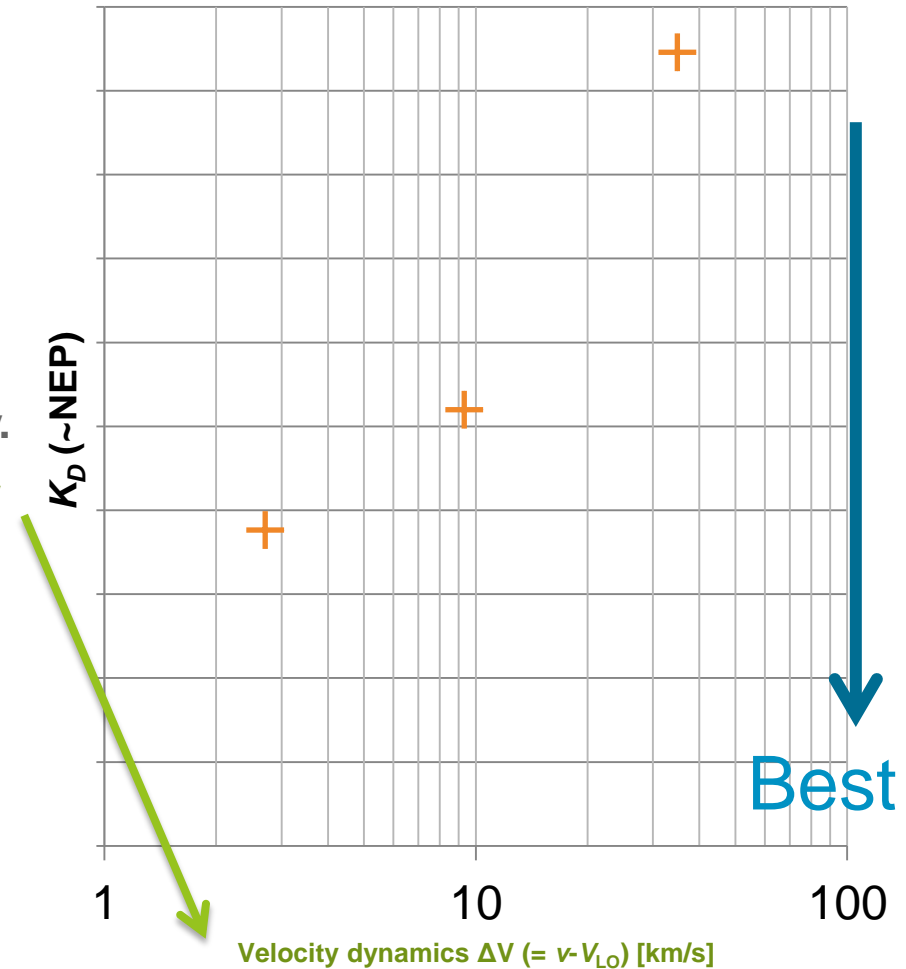
Dark noise

- Values taken from datasheet (NEP),
- Or measured.

$$\langle P_D \rangle = K_D(f_{BW}) \sqrt{f_{BW}}$$

- $K_D(f_{BW})$ is the NEP of the detector.
- It depends on the detector technology.
- In practice, it depends on the **detector bandwidth**.

NEP from a product line of detectors



WHICH TERMS ARE IMPORTANT FOR THE NOISE? 3/4

Digitizer noise

- The constant term (1) may be canceled with an offset
- The beating term (2) will set the dynamics of the digitizer, and its noise level.
- Mainly used characteristic to estimate digitizer noise : ENOB (effective number of bits), composed of:
 - noise floor level,
 - dynamic noise level,
 - distortions (total harmonic – TH).

$$K_Q = \frac{\Delta U_{P-P}}{2\sqrt{3} (2^{ENOB} - 1) f_{BW}}$$

- NB : distortions create harmonics, it is not so disadvantageous for free surface velocities



WHICH TERMS ARE IMPORTANT FOR THE NOISE? 4/4

ASE noise

- Generated by laser source, local oscillator and optical amplifier (EDFA).
- We do not use EDFA.
- We use very high-coherency laser (frequency width < 100 kHz), the spontaneous emission is really weak!
- For our PDV systems,

$$K_S = 0$$



TOTAL « NEP » K_T OF THE PDV SYSTEM IS:

$$K_T^2 = K_Q^2 + K_D^2 + K_{ph}^2 + \cancel{K_S^2}$$

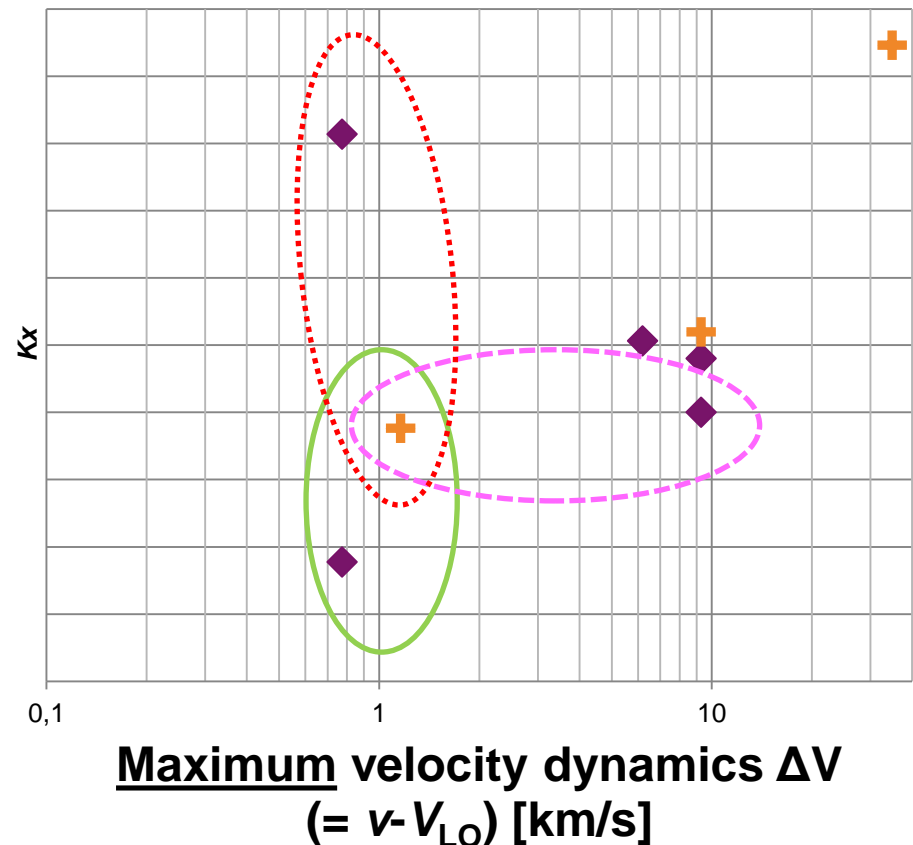
The digitizer may be changed

The detector may be changed

You can't act on it!
Changing P_{Lo} ? → [Backup](#)

How to select a couple of {detector + digitizer}?

◆ Digitizer + Detector



NOISE SOURCE LEVELS (8-GHZ CEA PDV SYSTEM)

Example with a 8-GHz bandwidth, CEA PDV system
($\alpha_{LO}P_{LO} = 250 \mu W$).

1. Photonic noise – $1.1 \mu W$

For low reflected power ($P_o < 10 \mu W$), it depends only on $\alpha_{LO}P_{LO}$ and detector efficiency.

2. Internal detector noise – $2.1 \mu W$

Depends on the detector.

3. Quantification noise – $2.3 \mu W$

Depends on the total reflected power and on the digitizer.

4. AES noise – $0 \mu W$

Very spectral-narrow lasers and no optical amplifier.

$$\text{So, } B_{det} = \sqrt{1.1^2 + 2.1^2 + 2.3^2} = 3.3 \mu W$$

[Curves of noise sources as a function of the vertical resolution \(backup\)](#)

WHY HETERODYNYING IS SO SENSITIVE ?

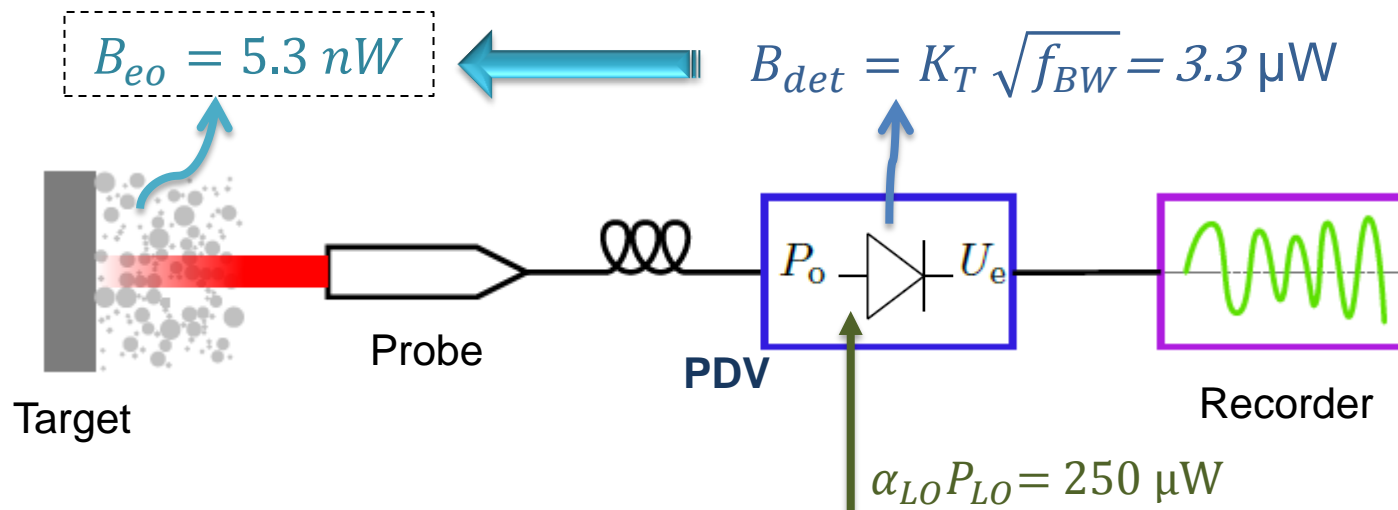
*Let's define a "NEP" in object space : "object-equivalent noise"
...(see my thesis)*

$$B_{oe} = \frac{K_T^2 f_{BW}}{8\alpha_o \alpha_{LO} P_{LO}} = 5.3 \text{ nW} \quad (\alpha_{LO} P_{LO} = 250 \text{ } \mu\text{W})$$

Another reinterpretation of the high detectivity of heterodyne systems:

$$B_{oe} \ll B_{det}, \text{ thanks to the } \sqrt{P_{LO} P_P} \text{ term.}$$

"Only objects, of which reflected and collected power is above B_{oe} , could be measured (using the raw signal)"

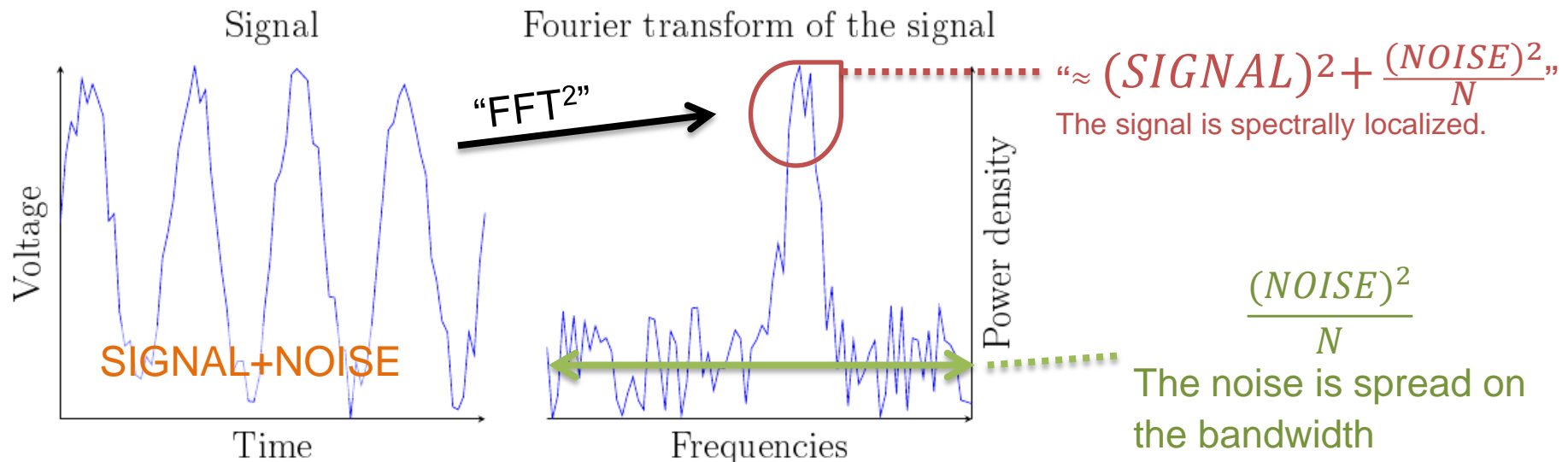


THE EFFECTS OF STFT TRANSFORM ON NOISE

WHY {HETERODYNING+STFT} ARE REALLY SO SENSITIVE?

• Fourier transform consequences:

- Mono frequency base: $e^{4i\pi \frac{v_m t}{\lambda_s}}$ -> monokinetic contributors are identified for each velocity (v_m).
- Noise is spread on all the bandwidth: after “ FFT^2 ”, noise mean level is reduced by N .



For example, see Harris, 1978, Proc IEEE

NOISE ON THE SPECTROGRAM: B_φ

Taking into account FT, noise level on the spectrogram is

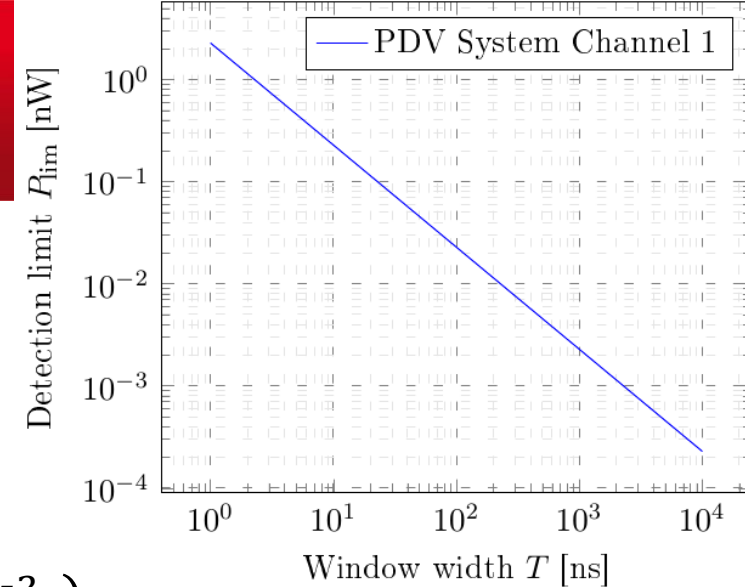
$$B_\varphi = \frac{1}{\hat{G}_i T f_{BW}} B_{eo}$$

$$= \frac{1}{\hat{G}_i T} \frac{1}{8\alpha_o \alpha_{LO} P_{LO}} \underbrace{(K_Q^2 + K_D^2 + K_{ph}^2)}_{\text{noise levels}}$$

\hat{G}_c is a window correction factor (≈ 0.2).

STFT makes detection limit independent of bandwidth (if Nyquist criteria is respected).

- but K_X are hardware-depending.
- It is useless to pre-filter the signal with a low-pass filter (maybe to reduce the number of samples).

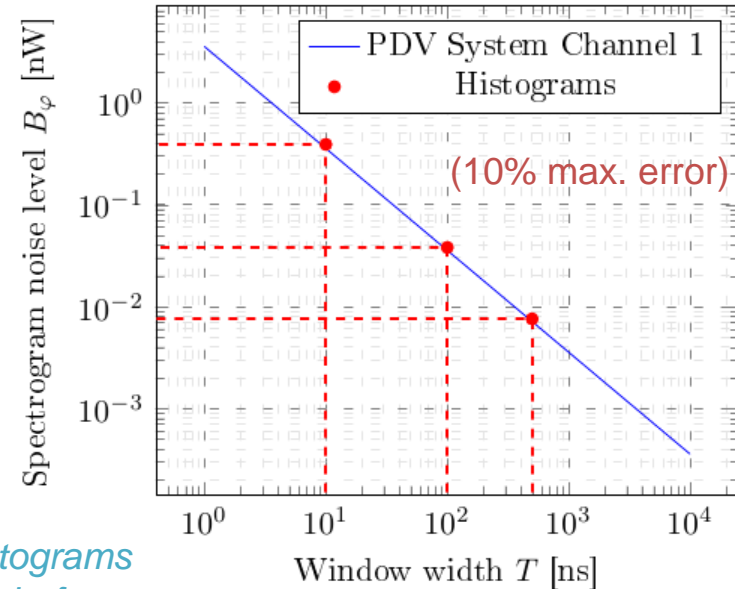


NOISE ON SPECTROGRAMS: COMPARISONS

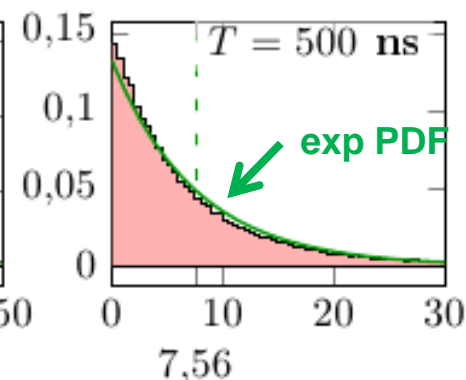
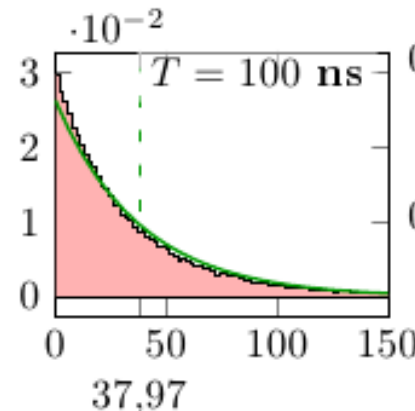
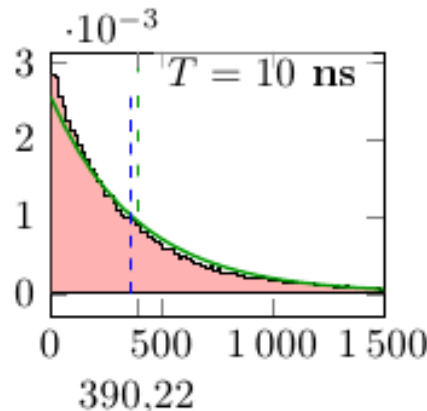
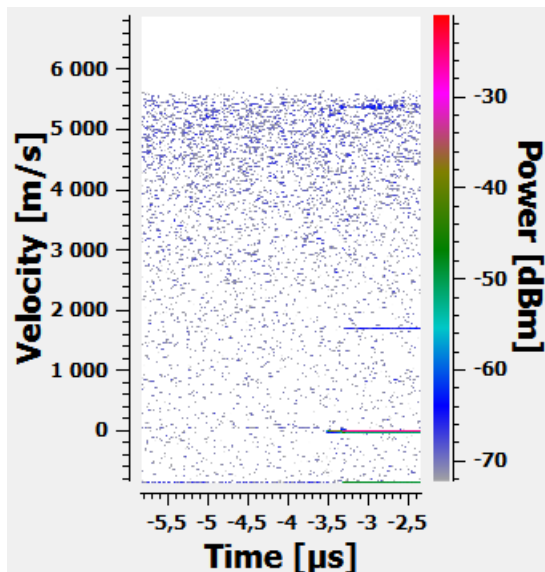
$$B_\varphi = \frac{1}{\hat{G}_i T f_{BW}} \quad B_{eo} = \frac{1}{\hat{G}_i T} \frac{1}{8\alpha_r \alpha_o P_r} \underbrace{(K_Q^2 + K_D^2 + K_{ph}^2)}_{\text{noise levels}}$$

- When no signal is present, **the noise PDF* is exponential**.
- B_φ is both the bias and the dispersion of the background noise (The mean is equal to the standard deviation).

*PDF: Probability Density Function



For 3 windows of T width, histograms computed from spectrograms before shock break-out are analyzed.



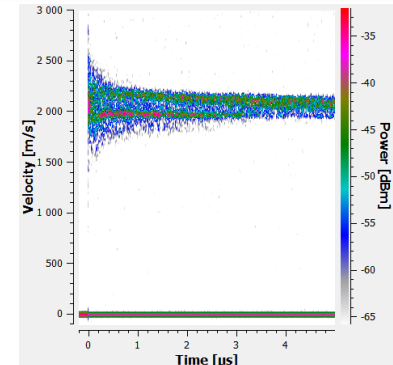
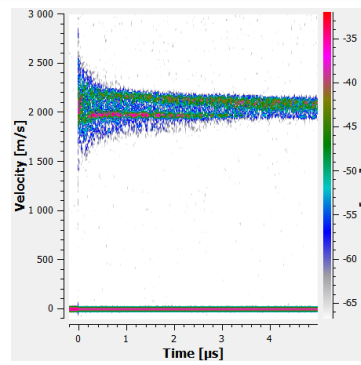
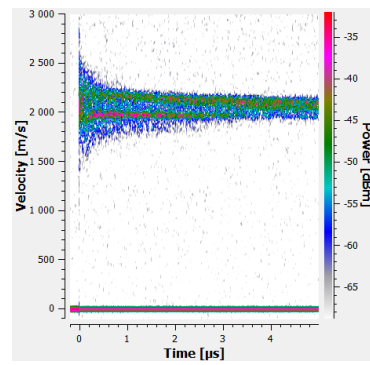
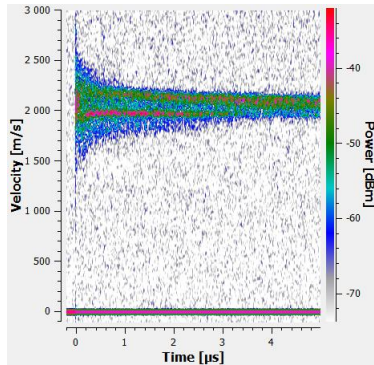
Power [nW]

Power [nW]

Power [nW]

HOW TO IMPROVE SPECTROGRAM DISPLAYS?

Which threshold should be set to display spectrogram?



How many “blue points” do we accept?

■ exp PDF

■ $z \rightarrow 1 - \exp\left(-\frac{z}{B_\varphi}\right)$ CDF : relative number of background values below z .

■ So, $s \rightarrow \exp\left(-\frac{s}{B_\varphi}\right)$ gives the relative number of « noisy » visible points.

■ Example : $T = 50 \text{ ns}$ @ 50 GS/s ,

■ $N = 2500$

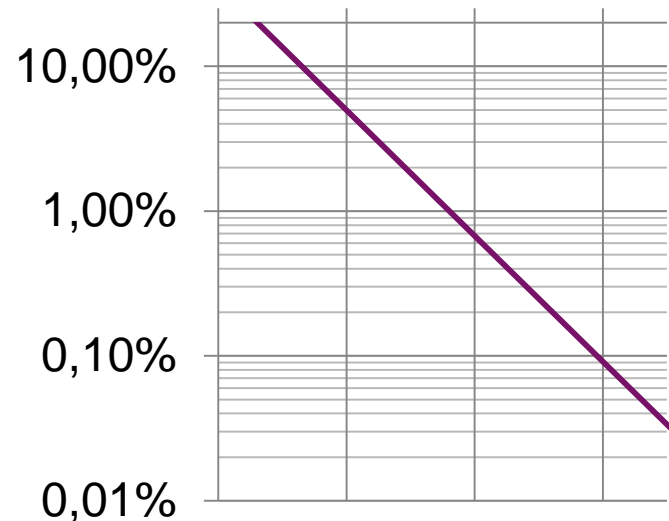
■ $S_\varphi = 5B_\varphi \rightarrow \approx 0.67 \%$

■ 17 expected noisy points in each column.

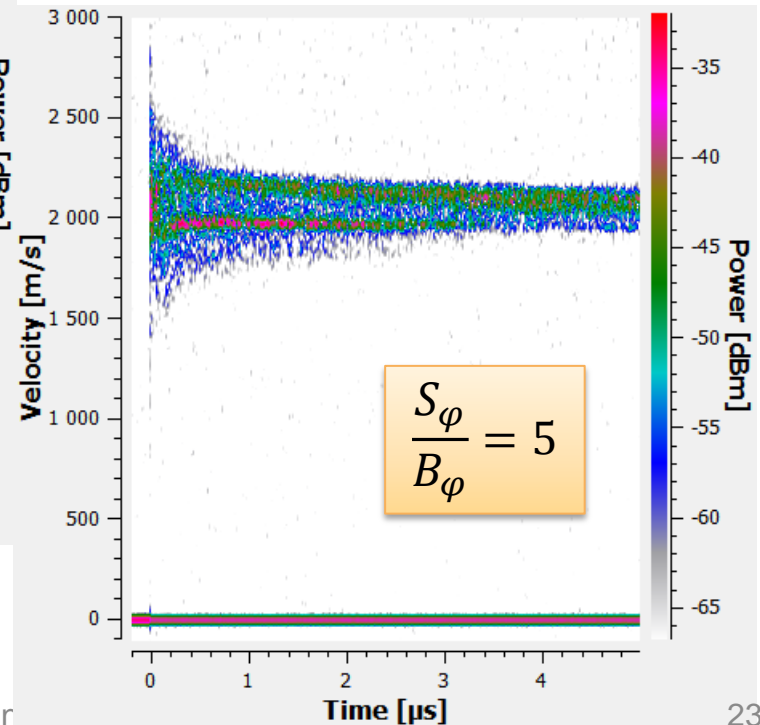
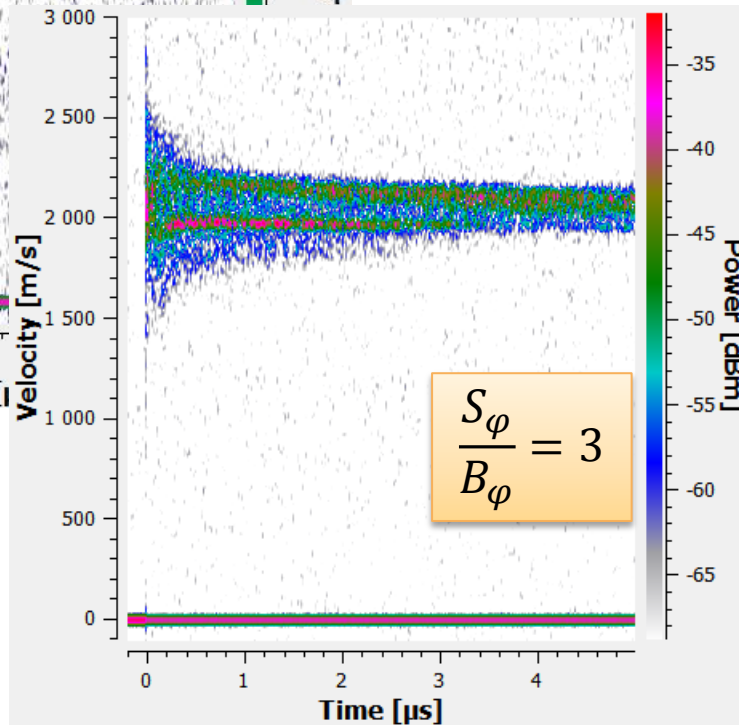
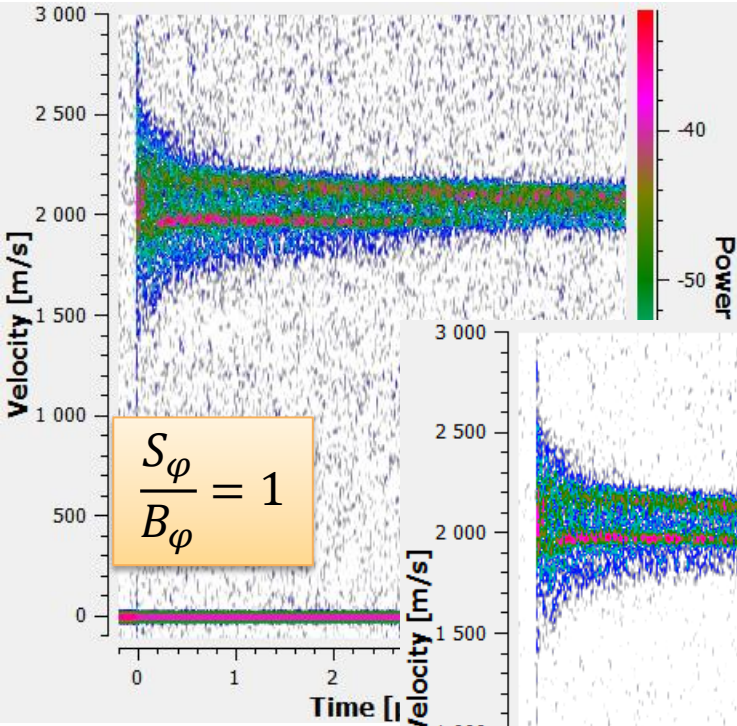
s : threshold

1B φ 3B φ 5B φ 7B φ

relative number of
« noisy » visible points



HOW TO IMPROVE SPECTROGRAM DISPLAYS? EXAMPLES



$\frac{S_\varphi}{B_\varphi}$	$\exp\left(-\frac{S_\varphi}{B_\varphi}\right)$
1	36 %
3	5,0 %
5	0,67 %
7	0,091 %

We present rules for PDV design in order to decrease noise:

- Contributors are quantified.
Major ones: detector background and digitizer noise.
Other source noises? (non-linear effect, backscattering inside the fiber, ...)
- Our model provides a good estimation of B_φ , the spectrogram mean noise level.

We analyzed the impact of noise on STFT spectrograms

- Noise level depends on noise density (“NEP”) and not on bandwidth,
 - Reducing digitizer bandwidth will not decrease B_φ (as long as Nyquist is respected).
- We are able to threshold the spectrogram as a function of the number of noisy samples.

THANK YOU FOR YOUR ATTENTION

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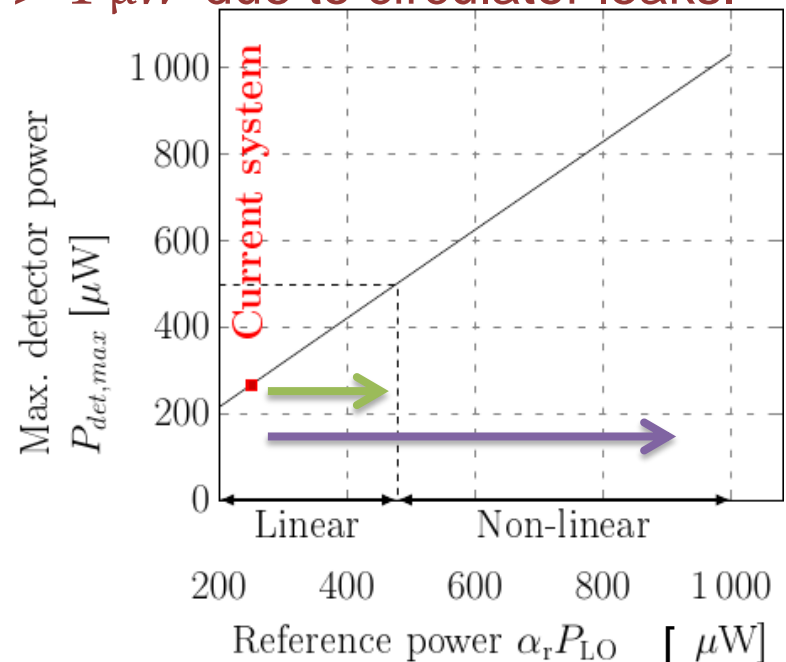
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BACKUP

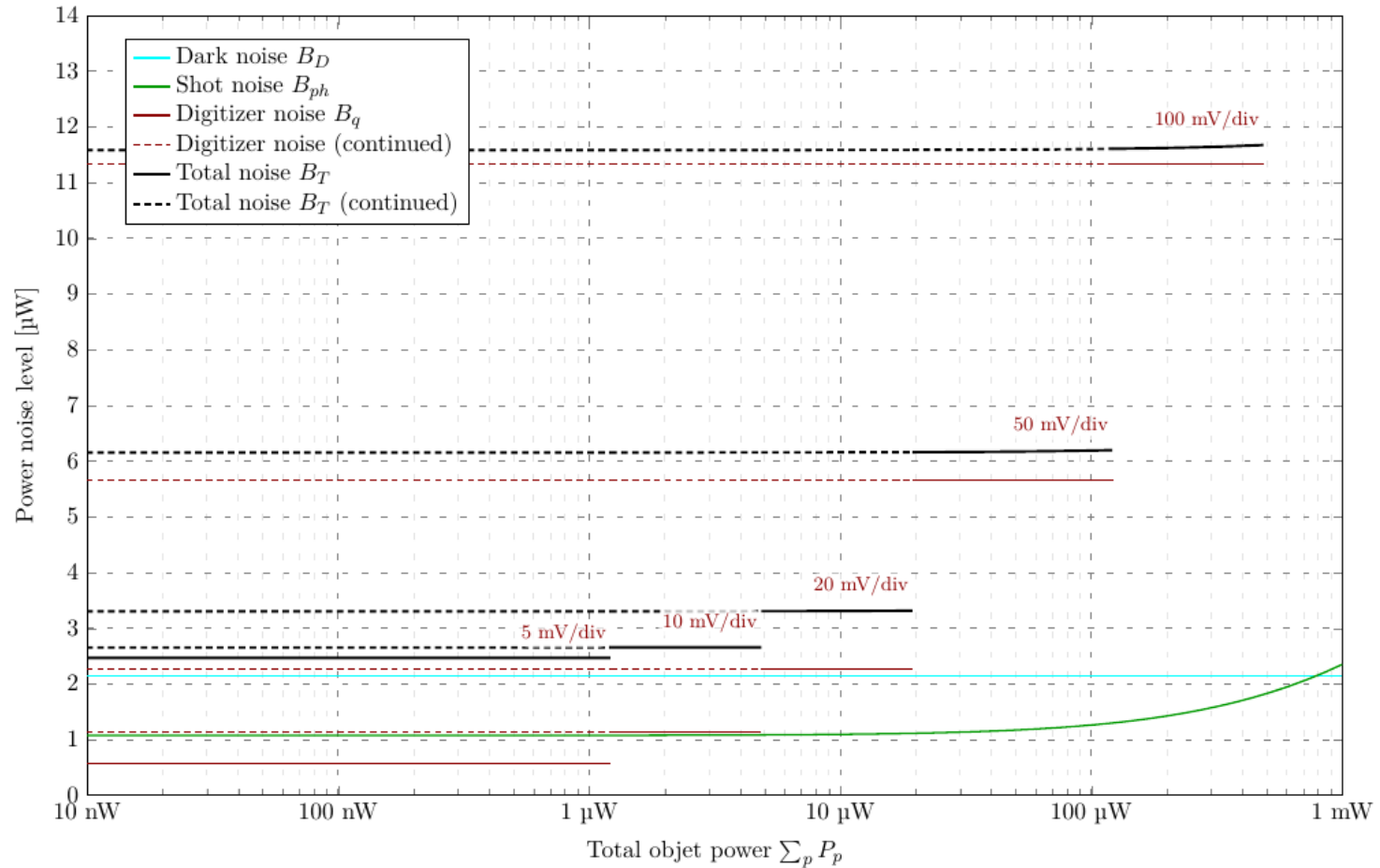
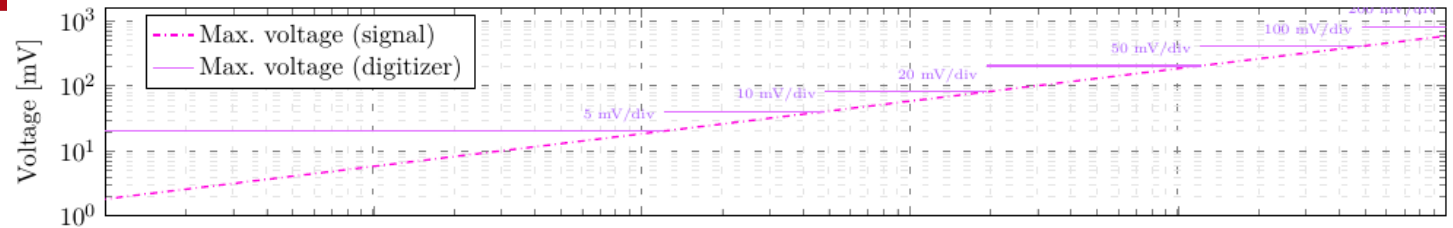
P_{LO} MAY BE INCREASED?

- Why? $P_{LO} \nearrow$, gain $\nearrow (\sqrt{P_{LO}})$, detectivity \nearrow because CEA PDV system is not photonic-noise limited.
- Detector limitation:
 - 1 mW for saturation.
 - $\geq 500 \mu\text{W}$ for non-linear effects, but may be corrected by post-processing (required P_{LO} or P_{det} knowledge).
- $P_{det,max} \approx \underbrace{\alpha_{LO}}_{\approx 2\%} P_{LO} + \sqrt{\alpha_{LO} \underbrace{\alpha_o}_{\approx 1} P_{LO} P_o}$, $P_o > 1 \mu\text{W}$ due to circulator leaks.
- Up to 4x P_r increasing is relevant ($K_{ph} = K_q$).
 - 2x: linear domain.
 - 4x: non-linear domain.

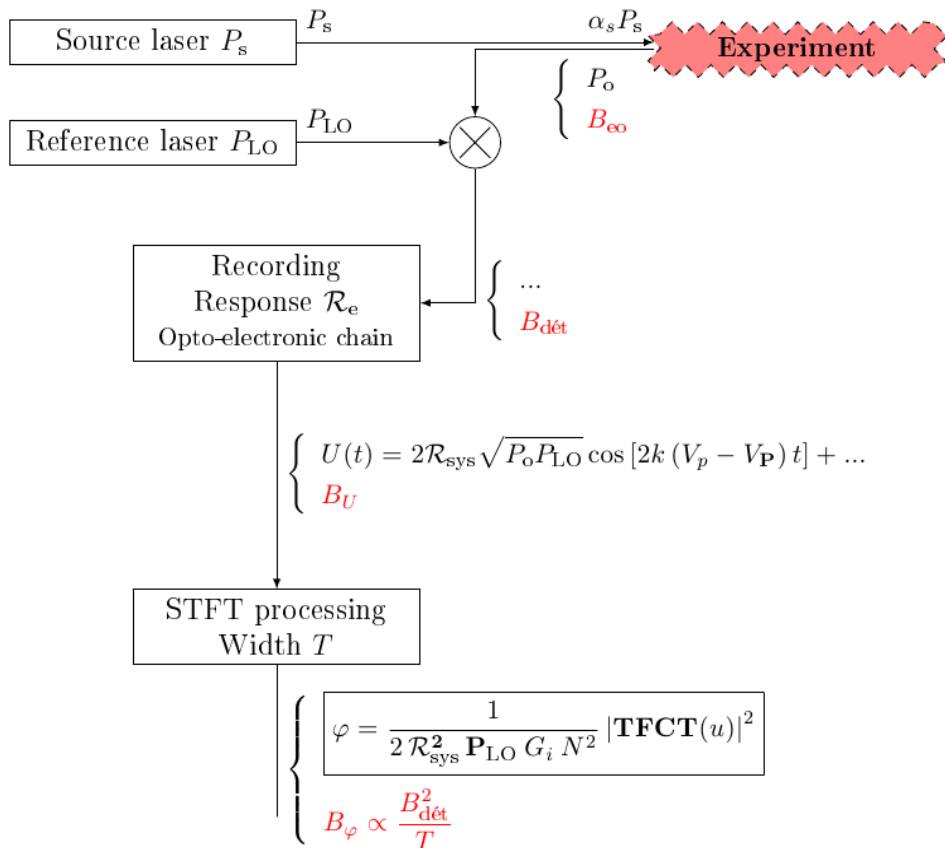
Technical development required.



NOISE AS A FUNCTION OF DIGITIZER V. RESOLUTION (CEA PDV SYSTEM)



“PDV RADIOMETRIC SCHEME”



Radiometric PDV equation of the spectrogram:

$$\varphi(t, v) = \frac{1}{2 \mathcal{R}_{sys}^2 P_{LO}} \frac{1}{\hat{G}_i N^2} \left| \text{STFT}_{f=2v/\lambda}^t(u) \right|^2$$

\mathcal{R}_{sys}^2 : system response,

P_{LO} : local oscillator power,

$\hat{G}_i N^2$: FT normalization, \approx bin number of the window (N).

➤Equals to the power distribution over the velocities.

➤Leads to the definition of a spectrogram mean noise level B_φ .

➤Do not take into account the polarization-efficiency ($\vec{e}_o \cdot \vec{e}_{LO}$) dot factor. Let's include it into \mathcal{R}_{sys} . (see J-E. FRANZKOWIAK's talk).

The point is:

The dB-expressed DSP is equal to the optical power distribution (with an offset).